

THE GROWTH OF SINGLE CRYSTAL OF ZINC SELENIDE

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ABSTRACT

ZnSe single crystal is made by the Bridgeman method under high inert gas pressure and its crystallographic characteristics are described. The lattice parameter of grown crystal is $5.667 \pm 0.003 \text{ \AA}$ and the twin structures along the $\langle 111 \rangle$ directions are observed. There are two kinds of $\{111\}$ planes in zinc blend type compounds, one is designated as $\{111\}$ surfaces or A surfaces and the other as $\{\bar{1}\bar{1}\bar{1}\}$ surfaces or B surfaces. Their differences are determined by means of the anomalous dispersion of X-ray and are discussed in relation to etch patterns.

§. 1. INTRODUCTION

As II-VI, III-V compounds (ZnS, ZnSe, CdS and GaP etc.) decompose into their volatile components far below their melting points, these single crystals have been grown mostly by chemical transport and vapour phase methods so far. The procedure of growing by these methods is very troublesome, taking a long time, and yet the grown crystal is small. However, if the vapour pressure is suppressed by high pressure, it is possible to grow the single crystal from the melt. A. G. Fisher has given a full detail of the crystal growth of the decomposable materials from the melt under pressure by variety methods¹⁾. The high pressure, high temperature electric furnace, developed by Gakei Electric works Co. Ltd., can be used as both Bridgeman and Tamman methods. The crystallization by means of the electric furnace of this type are reported by many investigators, such as ZnS by S. Shionoya²⁾, EuS by T. Nakayama *et al*³⁾, ZnSe by S. Ibuki *et al*⁴⁾ and S. Fukai *et al*⁵⁾.

The authors have succeeded in crystallizing ZnSe single crystal by Bridgeman method with good reproducibility. Advantages of Bridgeman method from the melt under high inert gas pressure are as follows :

- (1) The growth period is shorter than the other methods (chemical transport, growth from vapour phase).
- (2) The large single crystal can be grown.
- (3) Incorporation of impurities are easy.
- (4) Good reproducibility.
- (5) The homogenous crystal can be obtained.
- (6) The crystallization of solid solution is possible.

Before the measurement of optical and electrical properties, their surfaces are investigated. ZnSe has a zinc blend structure which is the same as that of diamond except for the alterna-

tion of the two different elements on successive lattice sites. Their space group is denoted as $F_{43m}(Td_2)$. The surfaces of compounds with zinc blend structure exhibit their own different characteristics, *i. e.*, the crystal growth, the etch pits formation and the gas adsorption as a result of the crystallographic polarity along the $\langle 111 \rangle$ directions. The surfaces $\{111\}$ terminating with group II atoms (Zn) are designated as the $\{111\}$ surfaces or the A-surfaces, and those terminating with group VI atoms (Se) are designated as the $\{\bar{1}\bar{1}\bar{1}\}$ surfaces or the B surfaces. The method of determination of polarity is reviewed by T. Ogawa⁶⁾ and the practical examples on III-V and II-VI compounds except for ZnSe can be seen in the reports by H. C. Gatos and E. P. Warekois et al.^{7,8)} The present paper describes the vertical Bridgeman growth method for ZnSe single crystals from the melt under pressure. In order to investigate the crystallographic characteristics of grown crystals, X-ray diffraction patterns, etching behaviours of $\{110\}$, $\{111\}$ surfaces and crystallographic polarity of the $\{111\}$ surfaces are studied. Furthermore, discussions are made on light figures, one of the conventional methods to determine the orientation of the crystal.

§. 2. CRYSTAL GROWING METHOD

The photograph of the electric furnace made by Gakei Electric Works Co. Ltd. is shown in Fig. 1. This furnace is filled with Ar gas and kept about 120 kg/cm² during crystal growth.

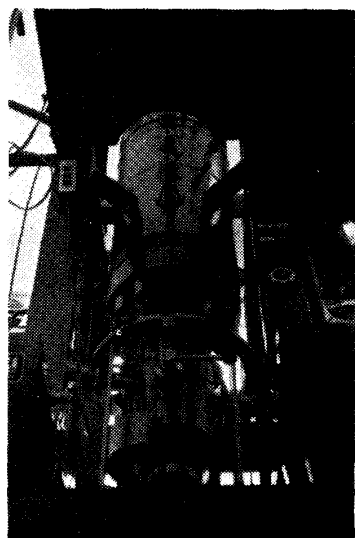


Fig. 1. Photograph of the high pressure, high temperature electric furnace.

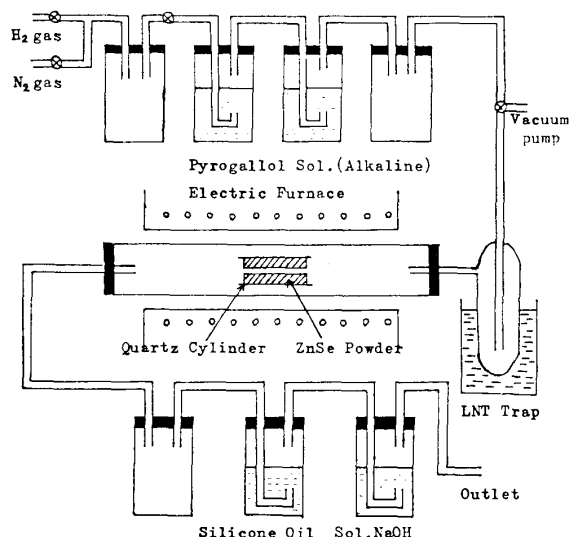


Fig. 2. Apparatus used for the refining of ZnSe powder.

The starting material (luminescent grade; 99.999 % purity; Sakai Chem. Co. Ltd.) is refined in the flowing H₂ gas at 1000°C for about 1 hour. Fig. 2. illustrates the apparatus used for this refining. The weight of the refined powder decreased to 80 % of the initial weight and the color of the powder turned from brown to bright yellow. By this refining the excess selenium, ZnO and H₂O are excluded quite well, and the powder is completely crystallized into zinc blend structure.

The crucible (made of graphite; shown in Fig. 3) which influences on the purity of the crystal must be carefully refined. The adopted procedures are as follows:

(1) Baked in the atmosphere of Ar gas at 1000°C for 1 hour.

(2) Immersed in the aqua-regia for 24 hours.

(3) Etched in the solution (100 cc HNO_3 +10cc HF +90cc H_2O) at 100°C for 12 hours.

(4) Rinsed in the distilled water at 70 - 80°C for 8 hours (repeated 4 times)

(5) Dried in the oven at 200°C for 24 hours.

(6) Baked in the vacuum at 1000°C for 1 hour.

The heating of high pressure furnace is carried out by P-I-D method with saturable core reactor. The crucible is maintained for 15-30 min. above the melting point ($1520^{\circ}\text{C}^{(9)}$) in order to be completely melted, then the crucible is pulled down from outside of the furnace by a magnetic coupling mechanism at an appropriate rate from about 1550°C to 800°C and the power is cut off. The temperature of the crucible is measured by thermocouple (W-5% Re/W-26% Re) which is attached to 18 mm under the end part of it. The temperature distribution in the furnace is shown in Fig. 4, and the temperature gradient calculated from Fig. 4. is shown in Fig. 5. These curves are

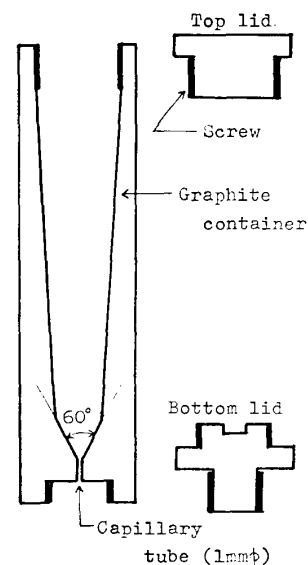


Fig. 3. Form of the graphite crucible.

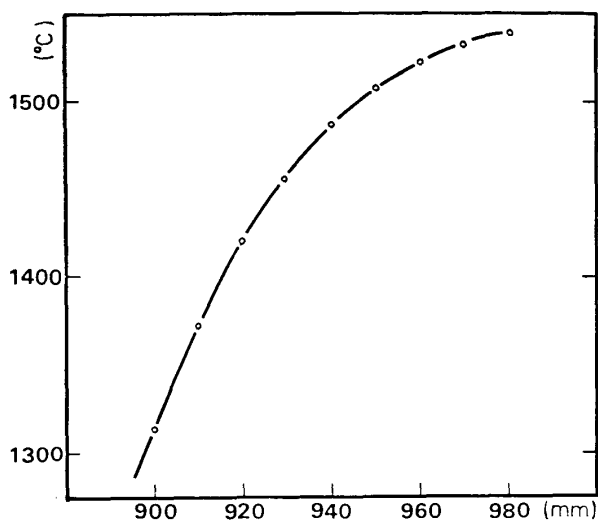


Fig. 4. Temperature distribution in the furnace measured by W-5% Re/W-26% Re thermocouple.

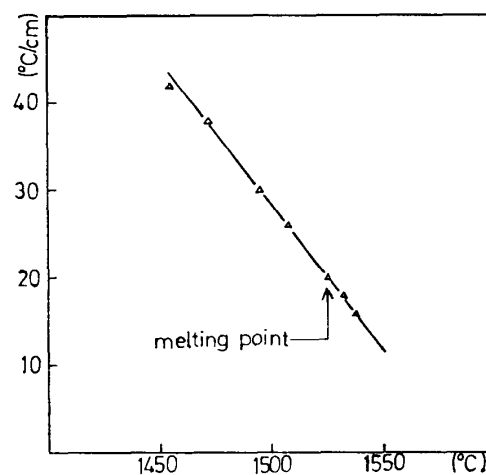


Fig. 5. Temperature gradient as a function of temperature. (calculated from Fig. 4)

uniquely determined by the heater shape and its surroundings, therefore the controllable factors which influence the crystal growth are the shape of crucible, gas pressure and the pulling down rate. In these factors the crucible is so influential on the crystallization that their inner wall is formed into the particular shape shown in Fig. 3. The capillary tube (1 mm ϕ) at the inner bottom promote the crystal growth of preferable orientation. By

Fig. 5., the temperature gradient near the melting point is about $20^{\circ}\text{C}/\text{cm}$, which is a little smaller than the reported value⁴⁾. Crucible temperature measured during crystal growth is shown in Fig. 6. Pulling down speed 18-36 mm/hour gives the single crystal without polycrystalline with good reproducibility. The evaluated temperature change per minute is represented in Fig. 7., which is all below $1^{\circ}\text{C}/\text{min}$. near the melting point. If the crucible

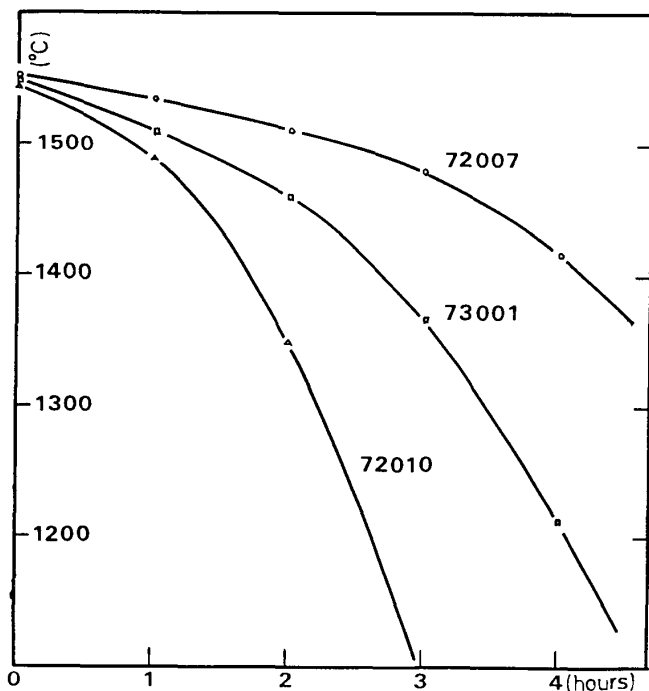


Fig. 6. Crucible temperature measured during crystal growth.

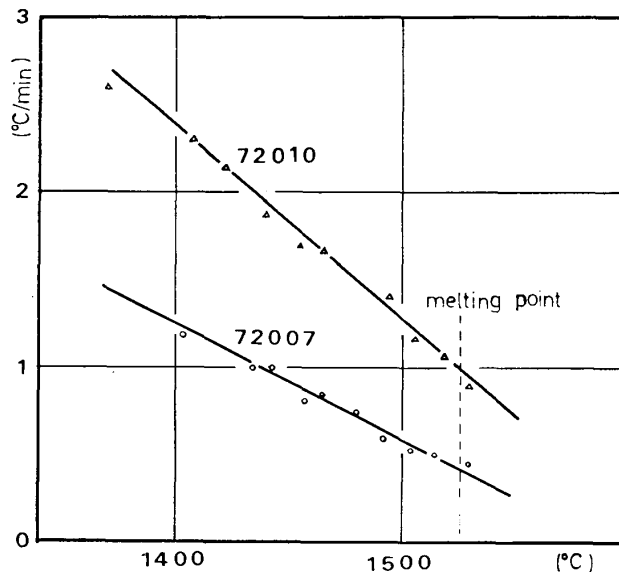


Fig. 7. Temperature change per minute of the crucible.

is pulled down at a speed below 18 mm/hour, the large single crystal is not grown since Bridgman method becomes similar to Tamman method and furthermore the deviation from the stoichiometry is caused in the long run. On the other hand, if the crucible is pulled down fastly (more than 35 mm/hour), the polycrystal is also grown since the equilibrium between the solid and liquid phases is not established. It is reported that the yield (grown crystal weight/charged weight) is less than 20 % and that grown crystal is polycrystalline under the pressure of $10 \text{ kg}/\text{cm}^2$ in the case of ZnSe ⁴⁾. The yield in our laboratory was about 90 % in the region from 90 to $120 \text{ kg}/\text{cm}^2$ and pressure had no remarkable effect on crystal growth. The direction of crystal growth was almost the $\langle 100 \rangle$, but the $\langle 110 \rangle$ and the $\langle 111 \rangle$ were observed sometimes. Generally speaking, the growth to $\langle 110 \rangle$ and $\langle 111 \rangle$ directions have the tendency to become polycrystalline. The obtained crystal is shown in Fig. 8., its size is about 17 mm in diameter and 65 mm in length. The as-grown crystal etched weakly in NaOH solution has the regular stripe patterns and the lapped crystal is orange-yellow in color except the end part and transparent. These stripes become disorder in polycrystalline ingot and are not observed in the worst condition. In such a case the crystal color is nearly blackish red. As-grown crystal has high resistivity (higher than $10^{10} \Omega\text{-cm}$) and slightly orange luminescence

is observed only at liquid nitrogen temperature by stimulating Hg 3650 Å line. The crystals annealed in the molten zinc (solvent extraction method¹⁰) show luminescence even at room temperature and their resistivity become few Ω -cm. The crystals used for the following experiment are all as-grown crystals.

§. 3. X-RAY ANALYSIS OF GROWN CRYSTAL

In order to investigate the grown crystal, X-ray diffraction patterns are examined about the raw material, refined powder and single crystal of ZnSe (Fig. 9). The raw material has the broad diffracted lines (Fig. 9-(a)), which means crystallization is not satisfactory. According to the qualitative fluorescent X-ray analysis, the raw material contains many impurities, such as K, Ca, Cl, Fe and S etc. In these impurities, S which would be incorporated in the process of synthesis of the ZnSe powder is predominant in quantity. The refined material show sharp diffraction lines of zinc blend type, and it is noted that they have a few hexagonal components, ZnO and other unknown materials. As far as X-ray analysis are concerned, the diffraction patterns of single crystal show no hexagonal components and oxides except for the end part of the crystal (Fig. 9-(c), (d)), which would be excluded during crystal growth. The evaluated lattice constant of ZnSe single crystal is $5.668 \pm 0.003 \text{ \AA}$, which is consistent with other reported results¹¹⁻¹³. The impurities contained in the single crystal are comparable to the refined powder by means of fluorescent X-ray analysis, so that the great care must

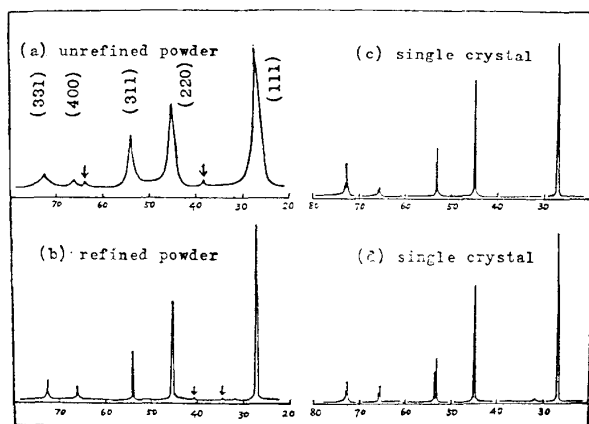


Fig. 9. X-ray diffraction patterns (Cu K_{α} radiation)

- (a) raw material (b) refined in the H_2 gas
(c) bottom part of single crystal
(d) upper part of single crystal.

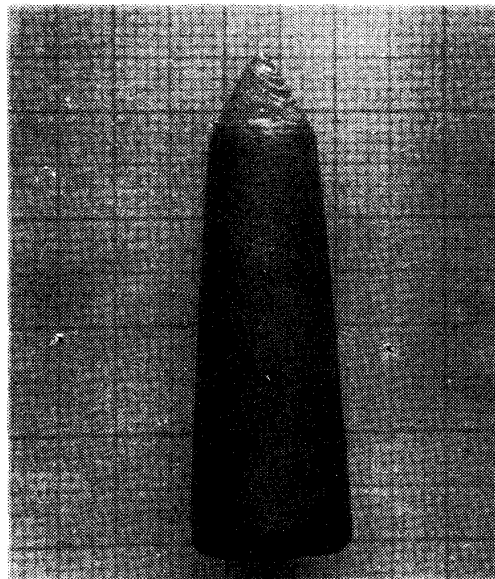


Fig. 8. Form of as-grown crystal of ZnSe (etched with NaOH solution).

be taken to the refining process of ZnSe powder and the cleaning of the crucible. It is very effective method for the exclusion of sulfur from the raw material to bake in the flowing H_2 gas. The back Laue photographs of the cleaved $\{110\}$ and the lapped $\{111\}$ planes are shown in Fig. 10. Both patterns do not show the patterns of single crystal. The pattern of the $\{110\}$ plane is 4-fold symmetry but not 2-fold symmetry and that of the $\{111\}$ plane is 6-fold but not 3-fold. These facts suggest that the crystal contains twins and their twin axis are the $\langle 111 \rangle$ directions. So the $\{110\}$ planes are composed of the (110) and $(\bar{1}\bar{1}0)$

planes. But if the $\{111\}$ planes are cut perpendicular to the cleaved $\{110\}$ planes, the plane without twin can be obtained. This is to be discussed later in detail.

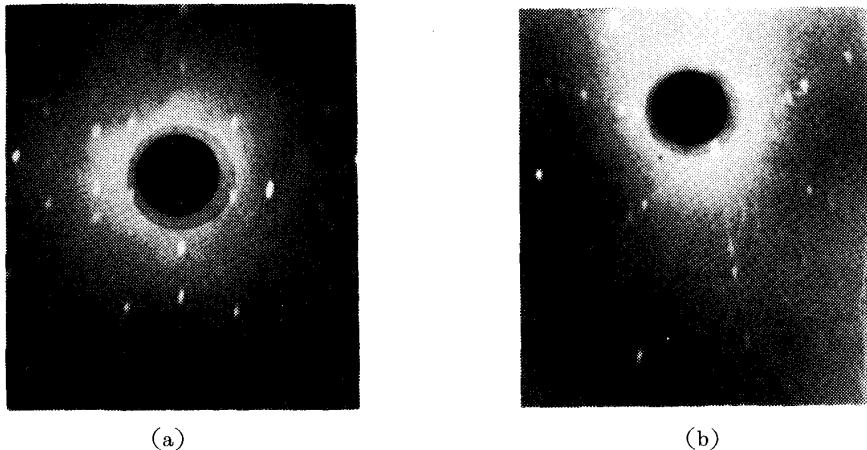


Fig. 10. Back Laue photographs of (a) $\{110\}$ and (b) $\{111\}$.

§. 4. ETCHING BEHAVIOURS OF ZnSe SINGLE CRYSTAL

(1) *Etch patterns and light figures of the $\{110\}$ surfaces.*

The etchant used in this experiment is the solution of 30 g NaOH+70 g H₂O. The other acid reagents, such as 1 (part by volume) conc HCl+1 conc HNO₃, HF *etc*, are examined, but NaOH reagent is the most suitable for the pits formation of ZnSe single crystal. The cleaved surface $\{110\}$ of ZnSe after etching at 110-120°C shows many small isosceles triangles as shown in Fig. 11. (a), (b), (c). Their base directions change at the parallel lines composed of high pits density on this plane, which are the crossing of the $\{110\}$ and $\{111\}$ planes. Referring to Laue photographs, the side direction of the triangles in etch pattern is $\langle 111 \rangle$

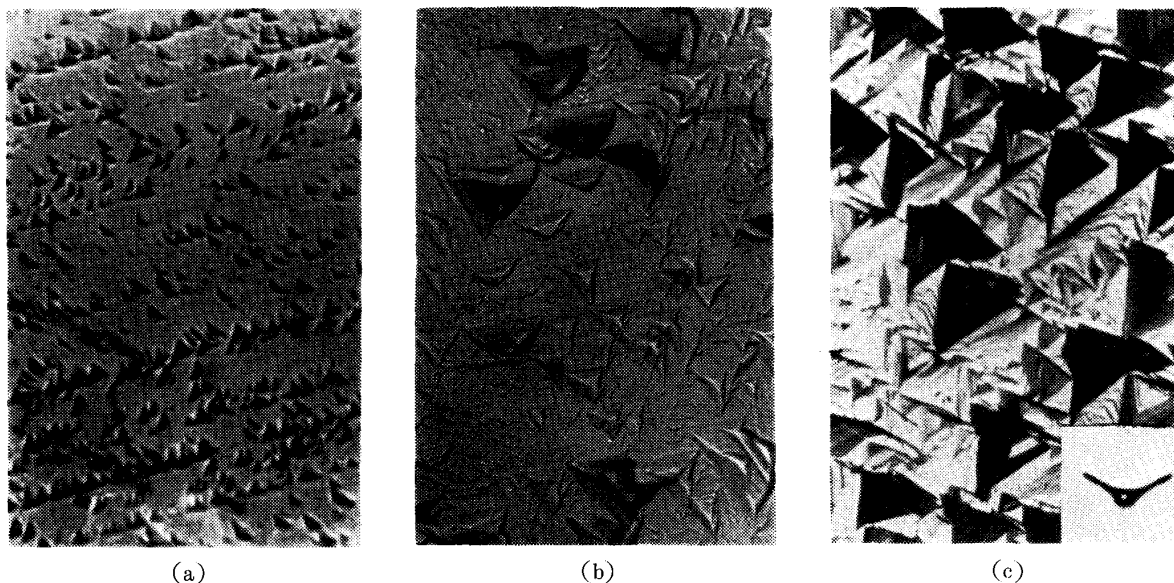


Fig. 11. Etch figures of $\{110\}$ (approx. $\times 400$); etching time (a) 2 min. (b) 5 min. (c) 10 min.

and the base direction is $\langle 110 \rangle$. The two types of triangular pits are not in the relation of the mirror reflection but the rotary reflection, that is, the etch pits are rotated 180° through the $\langle 111 \rangle$ twin axis.

As the observed angles of the triangular pits agree satisfactorily with calculated ones, if the pits orientations are determined with the X-ray photographs, the crystal orientation can be determined easily by the pits observation. When a cleaved faces are etched progressively, pits grow and combine at the twin boundary and their outlines become equilateral triangular pits. These pits growth can be explained by the difference of the chemical activity of the $\{111\}$ and $\{\bar{1}\bar{1}\bar{1}\}$ surfaces. To show the coincidence, their etch patterns of the pair of cleaved faces are shown in Fig. 12 (a), (b). Not only the twin boundaries but also the large pits might be caused by the inclusion matches closely on opposite side of cleaved surface. The pits density calculated from Fig. 11 (a) is about $1 \times 10^7/\text{cm}^2$ and from Fig. 11 (c) $2 \times 10^3/\text{cm}^2$. The twin density evaluated by the number of the twin boundaries per cm is between 600/cm and 1500/cm along the $\langle 111 \rangle$ directions. One of the determination methods for crystal orientation is to examine the light figure¹⁴⁾. The light figures are obtained from the $\{110\}$ surfaces etched heavily. The pattern shows two strong lines and a weak one. The weak line direction, that is, the bisector of the strong lines corresponds to the $\langle 111 \rangle$ directions (Fig. 11 (c)). For the slightly etched crystal, the light figure is not observed since the reflected beam is weak.

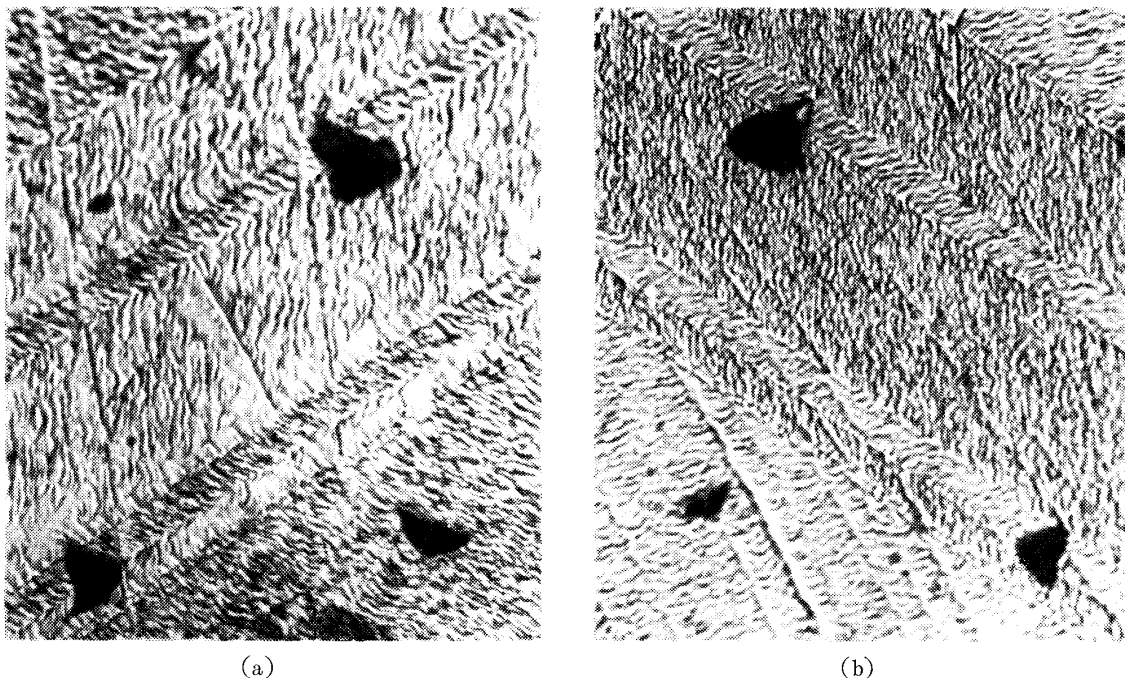


Fig. 12. Etch figures of pair of cleaved $\{110\}$ surfaces. (approx. $\times 400$)

(2) *Etch patterns and light figures of the $\{111\}$ surfaces.*

The $\{111\}$ surfaces of ZnSe after etched for 5 min. develop the triangular pits as shown in Fig. 13. There are two kinds of pits, sharp pits and diffuse ones. II-VI compounds

with zinc blend structure have the geometrical polarity along the $\langle 111 \rangle$ directions, i. e., the $\{111\}$ and $\{\bar{1}\bar{1}\bar{1}\}$ surfaces are nonequivalent to each other and exhibit the physical and chemical properties of their own.

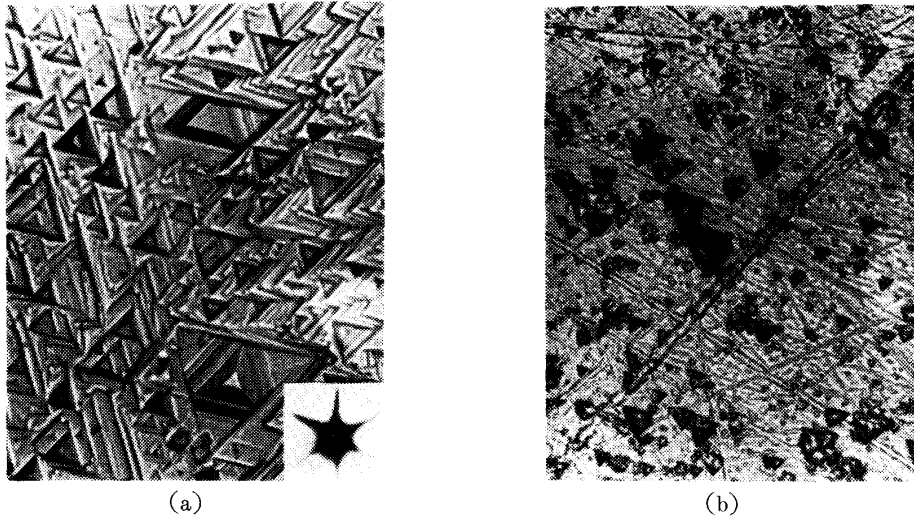


Fig. 13. Etch figures of (a) $\{111\}$ and (b) $\{\bar{1}\bar{1}\bar{1}\}$ surfaces. (approx. $\times 400$)

In zinc blend structure the atomic stacking sequences can be represented as $-A-B-C-A-B-C-$ etc. (Fig. 14¹⁵⁾). If the crystal contains the twins, this stacking is changed to $-A-B-\downarrow C-B-A-C-----B-A-B-C-A-B-C-$ etc.

The symbol \downarrow shows the composition surfaces. At the twin planes, the crystal rotates 180° through the $\langle 111 \rangle$ twin axis and the atomic arrangements of $-Zn-Se-Zn-Se-Zn-Se-$ etc. along the $\langle 111 \rangle$ directions are conserved.

So the distance between the first neighbours is equal.

If the crystal is cut parallel to this $\{111\}$ surfaces, the crystal of which surfaces have not twin is obtained.

The delta, nabra and overlaped pits appear clearly in Fig. 13 (a). These features can be explained by considering inaccuracy of cutting, thickness of the twinlamellae and the high etching rate. On the other hand, the etch pits of the Fig. 13 (b) (the opposite side shown in Fig. 13 (a)) have small delta pits only and the damages formed in the process of polishing are present yet. Etching rate of this surfaces is small, comparing with those of Fig. 13(a) for the same etchant. As the etching time

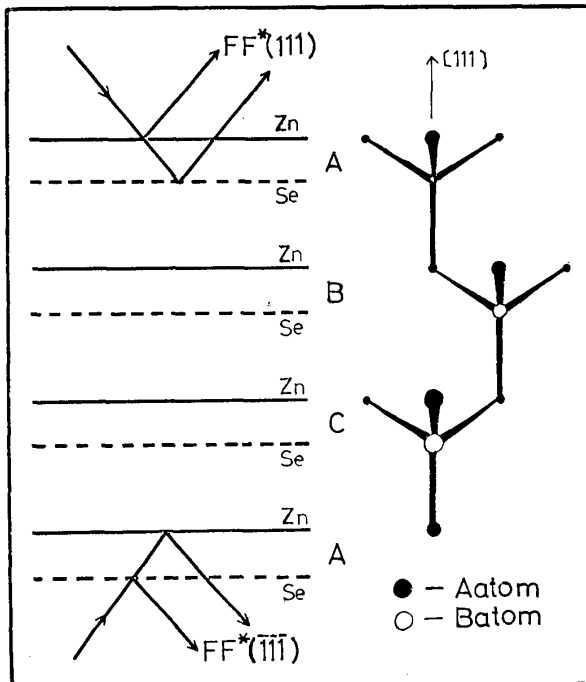


Fig. 14. Atomic stacking of zinc blend structure.

increases, the former surfaces (Fig. 13 (a)) become gradually mirror-like and then the flat bottomed pits disappear. Finally all remained pits become the shape of pyramid-like and their density decreases to about $5 \times 10^2/\text{cm}^2$. Whereas the latter surfaces (Fig. 13 (b)) remain rough and nonreflective, and pits or structure are recognized no longer (Fig. 15 (a), (b)). The surfaces etched with 1 (part by volume) conc. HCl+1 conc. HNO₃ show nearly the same characteristic patterns as those etched with NaOH etchant, that is, one is shallow rounded pits and the other a nondescript structure (Fig. 16). The fact that the pits are mainly produced only at one side of {111} is consistent with one that the dislocation pits are recognized in Zn plane in the cases of ZnS and ZnTe⁷⁾. In order to determine the crystallographic polarity of ZnSe, it is necessary to examine by means of anomalous dispersion of X-rays.

The light figure from the surfaces with sharp triangular pits is shown in Fig. 13 (a). This 6-fold symmetric pattern is expected from the etch pits shapes and consistent with the

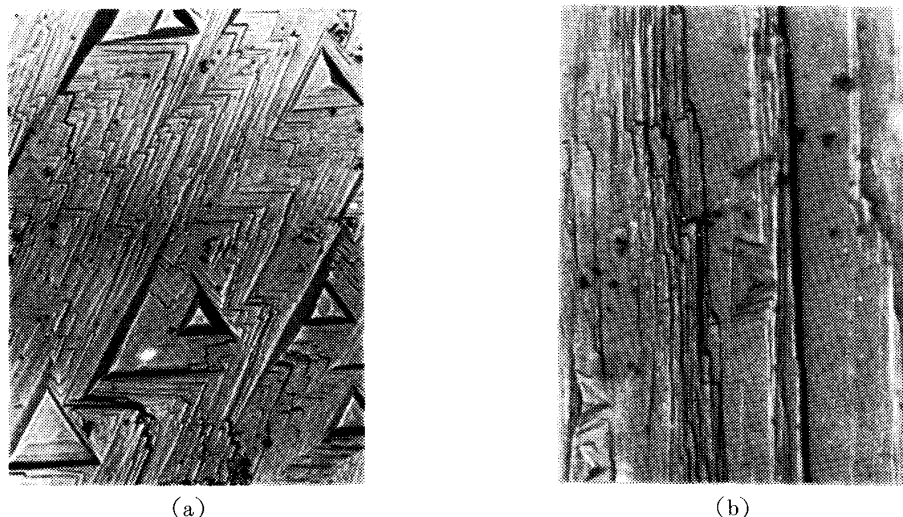


Fig. 15. Etch figures of the {111} Zn surfaces. (a) 10 min. etched (b) 20 min. etched. (approx. $\times 400$)

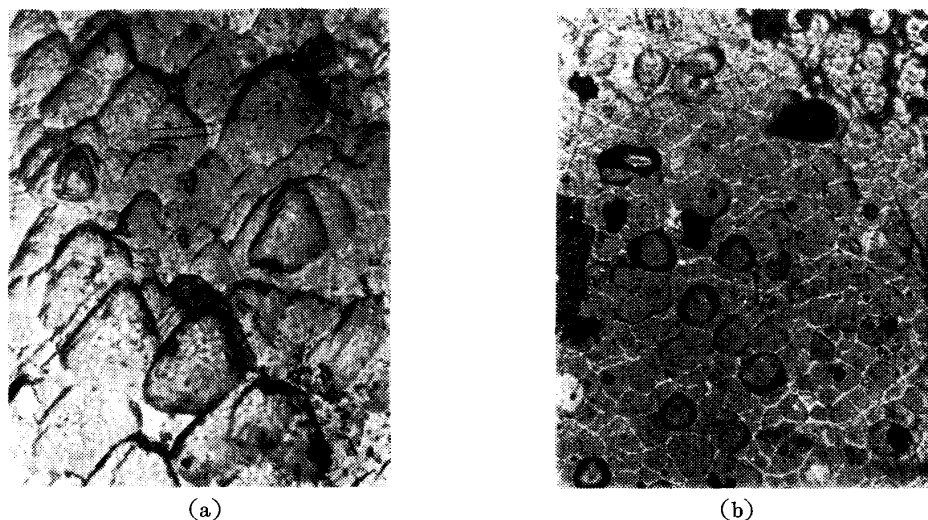


Fig. 16. Etch figures of the {111} surfaces produced with 1 (part by volume) conc. HCl+1 conc. HNO₃.
(a) {111} Zn surfaces (b) $\bar{1}\bar{1}\bar{1}$ Se surfaces. (approx. $\times 400$)

results of back Laue photographs. The light figure of the other $\{111\}$ surfaces can not be observed because of the etch pits diffuseness and shallowness.

§.5. DETERMINATION OF POLARITY

Recently it has been investigated that the identification of crystallographic polarity in II-VI compounds by employing X-ray diffraction and also correlated the etching characteristics of their A and B surfaces with X-ray identification by H. C. Gatos *et al*⁽⁸⁾. But they are not referred to ZnSe at all. In order to assign the polarity of ZnSe surfaces, the conventional fluorescent X-ray spectrometer (Rigaku; Geiger Flex SX) was used. As the analysing crystal the $\{111\}$ planes of ZnSe were used and KBr+RbCl, Ge single crystal and Cu metal were employed as characteristic X-ray sources. The schematic arrangements for the measurement are shown in Fig. 17.

The procedure of the calculation of the diffracted intensities are as follows,

$$F(hkl) = \sum_j f_j \exp 2\pi i(hx_j + ky_j + lz_j), \quad (1)$$

where f_j is the atomic scattering factors, h , k and l are Miller indices and x_j , y_j and z_j are the atomic partial coordinates in a unit cell.

$$\begin{array}{l} \text{Zn;} \quad 000, \quad \frac{1}{2} \frac{1}{2} 0, \quad \frac{1}{2} 0 \frac{1}{2}, \quad 0 \frac{1}{2} \frac{1}{2} \\ \text{Se;} \quad \frac{1}{4} \frac{1}{4} \frac{1}{4}, \quad \frac{4}{1} \frac{3}{4} \frac{3}{4}, \quad \frac{3}{4} \frac{1}{4} \frac{3}{4}, \quad \frac{3}{4} \frac{3}{4} \frac{1}{4}. \end{array}$$

Substituting these coordinates into eq. (1), the structure factor is obtained in the following equation.

$$F(hkl) = 4f_{Zn} + 4f_{Se} \exp(-in\pi/2). \quad (2)$$

The atomic scattering factors for Zn and Se layers considering the phase difference are given by

$$f_{Zn} = |f_{Zn}| \exp i\alpha_{Zn}, \quad f_{Se} = |f_{Se}| \exp i\alpha_{Se} \quad (3)$$

respectively. From eqs. (2) and (3) the following formula is obtained.

$$F(hkl) = 4 \{ |f_{Zn}| \exp i\alpha_{Zn} + |f_{Se}| \exp i(\alpha_{Se} - n\pi/2) \} \quad (4)$$

Therefore the structure factor is described as follows

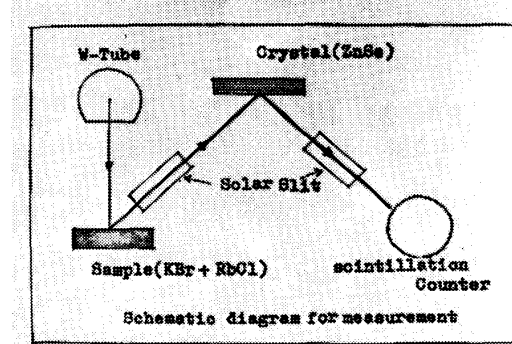


Fig. 17. Experimental set up for polarity measurement.

$$FF^*(hkl) = 16 \{ |f_{zn}|^2 + |f_{se}|^2 + 2|f_{zn}| |f_{se}| \cos(\alpha_{zn} - \alpha_{se} + n\pi/2) \} \quad (5)$$

The scattering factor corrected with dispersion by the K -electron is given by

$$f = f_o + \Delta f'_K + i\Delta f''_K \quad (6)$$

so that

$$|f| = \sqrt{(f_o + \Delta f'_K)^2 + (\Delta f''_K)^2} \quad (7)$$

This equation can be written quite nearly enough

$$|f| = f_o + \Delta f'_K + \frac{1}{2} \frac{(\Delta f''_K)^2}{f_o + \Delta f'_K} \quad (8)$$

The phase difference between the incident and the scattered radiation are derivated from eq. (6) as follows:

$$\tan \alpha_K = \frac{\Delta f''_K}{f_o + \Delta f'_K} \quad (9)$$

where f_o is the scattering factor for shorter waves than the wavelength of the absorption edge of the scattering element, that is, the scattering factor as usually tabulated. The values of $\Delta f'_K$ and $\Delta f''_K$, the real and imaginary parts of the correction to be applied the scattering factors on account of dispersion by K -electron, are calculated by means of Hönl's formulae.¹⁶⁾

$$\begin{aligned} \Delta f'_K &= \frac{2^7 e^{-4}}{9} \left\{ \frac{4}{x^2 (1 - \delta_K)^2} \cdot \log_e |x^2 - 1| - \frac{1}{(1 - \delta_K)^3} \left(\frac{2}{x^2} + \frac{1}{x^3} \log_e \left| \frac{x-1}{x+1} \right| \right) \right\}, \\ \Delta f''_K &= \frac{2^7 e^{-4}}{9} \pi \left\{ \frac{4}{x^2 (1 - \delta_K)^2} - \frac{1}{x^3 (1 - \delta_K)^3} \right\} \quad \text{if } x > 1, \\ &= 0 \quad \text{if } x < 1. \end{aligned} \quad (10)$$

where x is λ_{Kabs}/λ and λ being the wavelength of the radiation scattered, λ_{Kabs} the wavelength of the K absorption edge of the scattering element, and δ_K a parameter characteristic of the scattering element.

The values of δ_K are calculated for an element of atomic number Z from the formula

$$\delta_K = (A - 911/\lambda_{Kabs})/A, \quad (11)$$

where

$$\begin{aligned} A &= (Z - 0.3)^2 + 1.33 \times 10^{-5} (Z - 0.3)^4 + 3.55 \times 10^{-10} (Z - 0.3)^6 \\ &+ 11.7 \times 10^{-15} (Z - 0.3)^8 + \dots \end{aligned} \quad (12)$$

The values of $\Delta f'_K$ and $\Delta f''_K$ are tabulated by W. James¹⁶⁾, however, the authors calculated them directly from Hönl's formulae. The necessary parameters for the present calculation of

eqs. (8) and (9) are summarized in Table I.

Table. I. Basic parameters for the calculation of diffracted intensities.

Parameters {hkl}	$\frac{\sin \theta}{\lambda}$	f_0		A		δ_K	
		Zn	Se	Zn	Se	Zn	Se
(111) ($\bar{1}\bar{1}\bar{1}$)	0.153	25.56	28.67	892.4	1152.8	0.204	0.193
(333) ($\bar{3}\bar{3}\bar{3}$)	0.458	15.47	17.80				

Parameters Elements	$\lambda_{K\alpha}$ (Å)	$\lambda/\lambda_{K\alpha}$		$\Delta f_k'$		$\Delta f_k''$	
		Zn	Se	Zn	Se	Zn	Se
Rb	0.925	0.721	0.944	-0.252	-2.407	2.078	3.167
Br	1.041	0.811	1.063	-0.765	-2.962	2.536	0
Ge	1.255	0.978	1.281	-3.695	-1.912	3.428	0

From eq. (5), the diffracted intensities for the polar plane are given by

$$\begin{aligned}
 FF^*(111) &= 16 \{ |f_{zn}|^2 + |f_{se}|^2 - 2|f_{zn}| \cdot |f_{se}| \sin(\alpha_{zn} - \alpha_{se}) \}, \\
 FF^*(\bar{1}\bar{1}\bar{1}) &= 16 \{ |f_{zn}|^2 + |f_{se}|^2 + 2|f_{zn}| \cdot |f_{se}| \sin(\alpha_{zn} - \alpha_{se}) \}, \\
 FF^*(333) &= 16 \{ |f_{zn}|^2 + |f_{se}|^2 + 2|f_{zn}| \cdot |f_{se}| \sin(\alpha_{zn} - \alpha_{se}) \}, \\
 FF^*(\bar{3}\bar{3}\bar{3}) &= 16 \{ |f_{zn}|^2 + |f_{se}|^2 - 2|f_{zn}| \cdot |f_{se}| \sin(\alpha_{zn} - \alpha_{se}) \}.
 \end{aligned} \tag{13}$$

These factors can be calculated using the parameters listed in Table I. The results are shown in Table II.

Table. II. Calculated geometric scattering factors for ZnSe.

	$FF^*(111)$	$FF^*(\bar{1}\bar{1}\bar{1})$	$\frac{FF^*(111)}{FF^*(\bar{1}\bar{1}\bar{1})}$	$FF^*(333)$	$FF^*(\bar{3}\bar{3}\bar{3})$	$\frac{FF^*(333)}{FF^*(\bar{3}\bar{3}\bar{3})}$
Rb	1396	1293	1.079	441.2	504.6	0.874
Br	1151	1412	0.815	518.1	367.6	1.409
Ge	—	—	—	505.9	303.5	1.672

Actually, the values of $FF^*(hhh)/FF^*(\bar{h}\bar{h}\bar{h})$ are compared with the calculated ones. $CuK\alpha$ is used to correct the setting error of the crystal and to calibrate the surface condition of the both sides of the crystal, since anomalous dispersion does not appear. The typical patterns obtained from experiment are shown in Fig. 18.

The experimental results are summarized in Table III with the calculated values. Good agreement between observed and calculated values was obtained. If the surface aspect of the both sides of the crystal were equal, and a large effective area (it was $6 \times 10 \text{ mm}^2$ in

this experiment) were used, the more desirable results would be expected. Comparing the results in Table IV with etching patterns, it is decided that the plane with flat bottomed pits (Fig. 13 (a), Fig. 15 (a)) correspond to $(111)_{Zn}$ or the A surfaces and small rounded pits (Fig. 13 (b), Fig. 15(b)) to $(\bar{1}\bar{1}\bar{1})_{Se}$ or the B surfaces. This conclusion is consistent with the results obtained by K. F. Burr and J. Woods about the ZnSe crystal from the vapour phase.¹⁷⁾

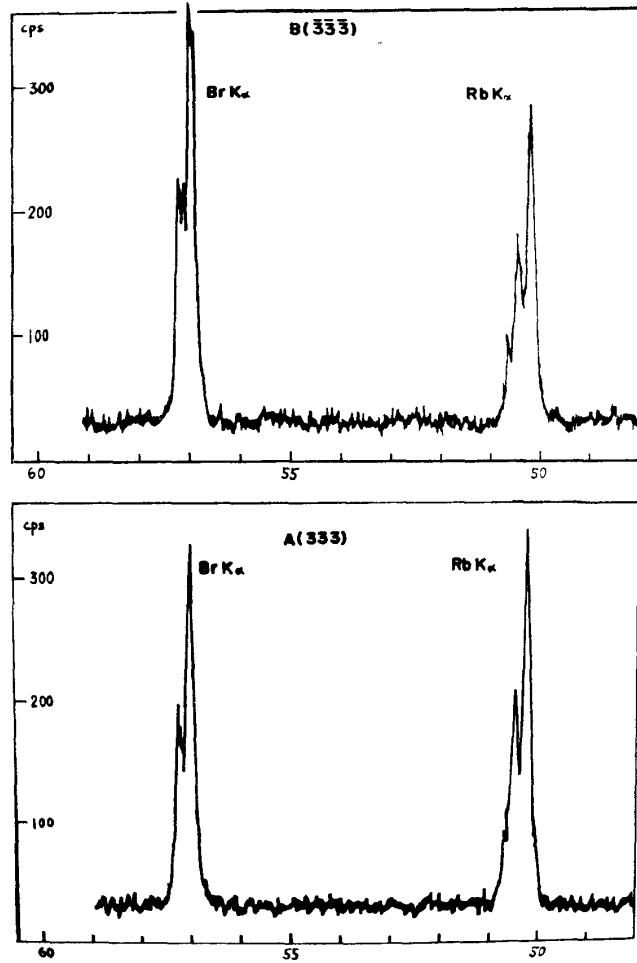


Fig. 18. X-ray $\{111\}$ reflection of ZnSe with fluorescence techniques.

Table. III. Geometric scattering factors (experimental).

$\{hkl\}$	Experimental			Calc.	Radiation
	$FF^*(hkl)$ cps	$FF^*(\bar{h}\bar{k}\bar{l})$ cps	$\frac{FF^*(hkl)}{FF^*(\bar{h}\bar{k}\bar{l})}$	$\frac{FF^*(hkl)}{FF^*(\bar{h}\bar{k}\bar{l})}$	
(333)	218	257	0.848	0.874	Rb- K_α
(333)	282	242	1.165	1.409	Br- K_α
(333)	740	470	1.574	1.672	Ge- K_α
(333)	1960	2180	0.900	1.000	Cu- K_α
(111)	950	620	1.532	1.079	Rb- K_α
(111)	890	860	1.035	0.815	Br- K_α

§. 9. SUMMARY

The single crystals of ZnSe were grown by Bridgeman method from the melt under high inert gas pressure. The grown crystals contain many twin and twinlamellae. The width in the adjacent lamellae could be estimated from the twin boundaries on the $\{110\}$ surfaces.

The evaluated width is about 30 microns along the $\langle 111 \rangle$ directions.

The etch patterns of the $\{110\}$ surfaces etched slightly show the small isosceles triangular pits clearly. When the surfaces are heavily etched, the equilateral triangular pits are developed clearly. The light figure of the $\{110\}$ surfaces is only observed from this heavily etched surfaces. The etch patterns of the $\{111\}$ surfaces are composed of the two kinds of equilateral triangles. The difference of these patterns can be explained by the chemical activity of the (111) and $(\bar{1}\bar{1}\bar{1})$ surfaces, *i. e.*, zinc and selenium layers. From the measurement of the ratio of diffracted intensities of X-rays $FF^*(hhh)/FF^*(\bar{h}\bar{h}\bar{h})$, the surfaces made of flat bottomed triangular pits correspond to $\{111\}_{zn}$ or the A surfaces and small rounded triangular pits to $\{111\}_{se}$ or the B surfaces. This explains the process from the isosceles to equilateral triangular pits on the $\{110\}$ surfaces.

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