

# Periodic Critical Orbits of Newton and Halley Iterations

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(Received November 1, 2000)

## 1. Introduction

In the previous paper [3] I have observed and conjectured that, for the 1-dimensional dynamics arising from iterations of certain typical 1-parameter families of smooth functions, parameter values at which periodic orbits of periods  $2^m(2n+1)$ ,  $n = 1, 2, 3, \dots$ , exist with multiplier  $-1$ , constitute geometrically convergent sequences.

In the present article we make a computer experiment for dynamical systems generated by Newton and Halley methods from polynomials of odd degree  $q$ ,  $q \geq 3$ . Unlike the previous case we have found parameter sequences at which periodic critical orbits of periods  $m, m+1, m+2, \dots$  (which forms a segment of the natural numbers) appear and, furthermore, that the corresponding asymptotic rates are *rational numbers* depending upon  $q$ .

In the case of these iterations I have observed that, on the left of main period-doubling cascade, two-way geometrically convergent parameter values admitting periodic critical orbits of odd periods and, in the right, reverse-ordered geometrically convergent parameters admitting super-stable periodic orbits of odd periods. This phenomenon seem to be similar to anti-monotonicity.

## 2. Fine Structures of Periodic Orbits

Let  $F(x)$  be a real-valued smooth function with a parameter  $\lambda$  and we consider one-dimensional dynamical system arising from iterating  $F$ . We shall denote generically by  $\lambda(p)$  a parameter value  $\lambda$  at which a super-stable periodic orbit of period  $p$  appears. In other words  $F(x)$  with  $\lambda = \lambda(p)$  admits a  $p$ -periodic point whose orbit contains a critical

point of  $F(x)$ . If, near to  $\lambda(p)$ , there is the parameter value  $\lambda$  at which a  $p$ -periodic orbit with multiplier  $-1$  appears, then we denote it by  $\omega(p)$ . Set  $\delta(p) = \omega(p) - \lambda(p)$ . We also consider the *diameter* of the periodic orbit of period  $p$  corresponding to  $\lambda = \lambda(p)$  and denote it by  $\delta \text{ orb}(p)$ , namely,

$$\delta \text{ orb}(p) = \text{the maximum of } p\text{-periodic points} - \text{the minimum of } p\text{-periodic points}$$

It occurs often that, if one takes  $p$ -periodic orbits for  $p = m, m+1, m+2, \dots$  in a suitable family, then the corresponding sequence  $\lambda(m), \lambda(m+1), \lambda(m+2), \dots$  becomes monotone. In such cases we denote the successive ratios of its first difference by  $\rho(p)$ . Namely

$$\rho(p) = [\lambda(p+2) - \lambda(p+1)] / [\lambda(p+1) - \lambda(p)].$$

Consider the case where the sequence  $\rho(p)$  has a limit  $\rho(\infty)$ , called the *asymptotic ratio*, as  $p$  tends to infinity. If  $\rho(\infty)$  is smaller than 1 then  $\lambda(p)$  converges to a limit  $\lambda(\infty)$ : if  $\rho(\infty)$  is larger than 1 then  $\lambda(p)$  diverges, in which we call  $\{\lambda(n)\}$  *geometrically divergent*.

Incidentally, it will often be observed that  $\{\delta \text{ orb}(p)\}$  for  $\lambda = \lambda(p)$  converges geometrically as  $p$  tends to infinity in which we denote the asymptotic ratio by  $\rho_\delta$ . Further it is often observed that the sequence  $\{\rho(p)\}$  is also geometrically convergent with the same asymptotic ratio as in  $\lambda(p)$ ; in such a case we denote the ratio by  $\rho_\rho$ .

If there is a geometrically convergent sequence  $\lambda(2n+1)$ ,  $n \geq m$ , thus  $\omega(0, n) = \lambda(2n+1)$  in the notation of [3], then we have similar ratios  $\rho(0, n)$  with asymptotic ratio  $\rho(0, \infty)$ .

In order to describe and classify parameter sequences appearing in 1-dimensional dynamics, we shall use the following notations. For a monotone parameter sequence

$$\lambda(m) < \lambda(m+1) < \lambda(m+2) < \dots \quad \text{or} \quad \lambda(m) > \lambda(m+1) > \lambda(m+2) > \dots$$

converging to  $\lambda(\infty)$  with asymptotic ratio  $\rho$ , we denote it by  $\text{NS}(m, \lambda(m), \rho; \lambda(\infty))$ .

If there is a sequence

$$\lambda(m) < \lambda(m+2) < \lambda(m+4) < \dots$$

with asymptotic ratio  $\rho$  and limit  $\lambda(\infty)$  then we denote it by  $\text{RS}(m, \lambda(m), \rho; \lambda(\infty))$ .

Similarly,  $\text{LS}(m, \lambda(m), \rho; \lambda(\infty))$  denotes the sequence

$$\dots < \lambda(m+4) < \lambda(m+2) < \lambda(m)$$

with asymptotic ratio  $\rho$  and limit  $\lambda(\infty)$ . The two-way sequence

$$\dots < \lambda'(m+4) < \lambda'(m+2) < \lambda(m) < \lambda(m+2) < \lambda(m+4) < \dots$$

with left and right asymptotic ratios  $\rho'$  and  $\rho$  will be denoted by  $\text{TS}(m, \lambda(m); \rho', \rho; \lambda'(\infty), \lambda(\infty))$ .

### 3. Newton Method

Applying Newton's method to a given smooth function  $f(x)$  we have

$$F(x) = Nf(x) = x - f(x)/f'(x) \tag{3.1}$$

In the frame-work of complex dynamics, iterations of  $F(x)$  have been studied extensively by many people ( see [1, 2, 4, 5 ] ).

The derivative of  $Nf(x)$  is given by

$$(Nf)'(x) = [ f(x) f''(x) ] / ( f'(x) )^2 \tag{3.2}$$

Hence any root of  $f'(x) = 0$  which is not a root of  $f(x) = 0$  gives a *free* critical point. In the light of theorems of Fatou or Singer, we usually start with such free critical points in case we search for (attractive)  $p$ -periodic points of  $Nf(x)$ .

Take for  $f(x)$  polynomials of *odd* degree  $q$  :

$$f(x) = x^q + (\lambda - 1)x - \lambda$$

Then  $x = 0$  is the free critical point. We have made computer-search of periodic points and corresponding  $\lambda(p)$  using a program similar to the one stated in [3]. We have found the following parameter-sequences:

for  $q = 3$  RS(7, 0.254338438687, 0.25141; 0.254619273771199) with  $\rho_\delta = 0.251413$   
 TS(3, 0.26744849; 0.251879, 0.259259; 0.256078731774102, 0.275682203651)  
 with  $\rho_\delta' = 0.25188, \rho_\delta = 0.25926$

NS(5, 0.398756368, 2/3; 0.4459873529729)

NS(3, 0.510650459, 2/3; 1) with  $\rho_\delta = 1.5$

for  $q = 5$  NS(5, 0.3406616321, 0.8; 0.34777632228) with  $\rho_\delta = 0.85$

NS(3, 0.3367402080, 0.8; 0.373062) with  $\rho_\delta = 0.8$

NS(3, 0.463469485, 0.8; 1) with  $\rho_\delta = 1.25$

for  $q = 7$  TS(7, 0.36607039; 0.250885, 0.2518235; 0.3658058158, 0.3667288329)

TS(5, 0.36847724; 0.2524546, 0.254468; 0.3673330067, 0.3691757001)

NS(3, 0.37218037, 6/7; 0.39910233) with  $\rho_\delta = 0.85$

NS(5, 0.423125779, 6/7; 0.448767458756) with  $\rho_\delta = 0.855$

NS(3, 0.45820702, 6/7; 1) with  $\rho_\delta = 7/6$

For even  $q$  we have never found periodic points except fixed points.

From these numerical evidences, it is reasonable to state :

**Conjecture N** : For odd  $q$  there are geometrically convergent sequences

$$\lambda(m), \lambda(m+1), \lambda(m+2), \dots$$

with asymptotic rate  $(q-1)/q$  whose orbits have diameters such that  $\rho_\delta$  are equal to  $(q-1)/q$  or  $q/(q-1)$

#### 4. Halley's method

Given a smooth function  $f(x)$ , the *Halley iteration function*  $Hf(x)$  is a special case of the König iteration function of  $f(x)$  (see Vrscay-Gilbert [5]) and defined as follows :

$$Hf(x) = x - (3-1) [ 1/f(x) ]' / [ 1/f(x) ]''$$

$$\begin{aligned}
&= x - f(x) / [f'(x) - 0.5 f(x)f''(x)/f'(x)] \\
&= x - 2 f(x)f'(x) / [2 (f'(x))^2 - f(x)f''(x)] \quad (4.1)
\end{aligned}$$

whose derivative is given by

$$\begin{aligned}
(Hf)'(x) &= [f(x) / \{2 (f'(x))^2 - f(x)f''(x)\}]^2 [3 f''(x)^2 - 2f'(x)f'''(x)] \\
&= -2 [f(x) / \{2 (f'(x))^2 - f(x)f''(x)\}]^2 SD(f(x)) \quad (4.2)
\end{aligned}$$

where  $SD(f(x))$  denotes the Schwarzian derivative of  $f(x)$  (see Yao [6]).

Take for  $f(x)$  polynomials of degree  $q$ :

$$f(x) = x^q + (\lambda - 1)x - \lambda$$

Then a simple calculation shows that

$$3f''(x)^2 - 2f'(x)f'''(x) = q(q-1)x^{q-3} \{q(q+1)x^{q-1} - 2(q-2)(\lambda-1)\}$$

Thus free critical points of  $Hf(x)$  are 0 and

$$\begin{aligned}
&e [2(q-2)(\lambda-1)q^{-1}(q+1)^{-1}]^{(1/(q-1))}, \text{ where } e = \pm 1 \quad \text{for odd } q, \lambda \geq 1 \\
&-[2(q-2)(1-\lambda)q^{-1}(q+1)^{-1}]^{(1/(q-1))} \quad \text{for even } q, \lambda < 0
\end{aligned}$$

We have made numerical computation for  $q = 3, 5, 7, 9, 11$  using a program similar to the one in [3], and obtained the following parameter-sequences:

for  $q = 3$  NS(5, 1.5865404026, 0.5; 1.6417742853)

NS(3, 1.92678793818, 0.5; 2.5) with  $\rho_\delta = 0.5$

for  $q = 5$  RS(3, 1.058514888, 0.174847384; 1.078447789469) with  $\rho_\delta = 0.1748474$

NS(5, 1.69885664, 2/3; 1.8764003026) with  $\rho_\delta = 2/3$

NS(3, 2.61686822, 2/3; 16.64114978425844) with  $\rho_\delta = 2/3$

for  $q = 7$  NS(5, 2.05839513, 0.75; 2.5782312276669) with  $\rho_\delta = 4/3$

NS(3, 5.85146057, 0.75; 220649.3075219) with  $\rho_\delta = 4/3$

TS(3, 257120.242413; 0.044540715, 0.0432574; 255141.185382, 259189.938888)

for  $q = 9$  NS(3, 1.060902495, 0.8; 1.002669187889) with  $\rho_\delta = 0.8$

RS(5, 1.0815720369, 0.1810494915; 1.086393470229) with  $\rho_\delta = 0.181$

NS(5, 2.828392423, 0.8; 4.822331861043367) with  $\rho_\delta = 0.8$

NS(7, 182006.8688802, 0.8; 177342.0166683) with  $\rho_\delta = 1.25$

NS(5, 641681.236982, 0.8; 602001.8463) with  $\rho_\delta = 1.25$

NS(9, 2005524.47254, 0.8; 95564064.938) with  $\rho_\delta = 0.8$

NS(9, 485762030372458522287635945.3, 0.8; 22269277604254934850600000)

with  $\rho_\delta = 0.8$

for  $q = 11$  RS(5, 1.00136823439, 0.192062126; 1.002195290795) with  $\rho_\delta = 0.1743$

LS(5, 1.06040025182, 0.178954; 1.05992846044719) with  $\rho_\delta = 0.017895$

RS(5, 1.0195165052, 0.1742719; 1.2213098025)

NS(5, 5.023110591, 5/6; 24.870979) with  $\rho_\delta = 1.2$

NS(6, 169128.315, 5/6; 2116878.842979) with  $\rho_\delta = 5/6$

- NS(5, 396189822567385.2, 5/6; 1123168.90337) with  $\rho_\delta = 5/6$   
 NS(9, 780523919991.023, 5/6; 76218176274791.62945) with  $\rho_\delta = 5/6$   
 NS(9, 21570955, 5/6; 259305220704360970) with  $\rho_\delta = 1.2$   
 $q = 6$  NS(5, -12.780262, 5/7; -412.001234) with  $\rho_\delta = 7/5$   
 NS(5, -2.2974041, 5/7; -0.477889078259) with  $\rho_\delta = 7/5$   
 We could not find periodic points of period 3.  
 $q = 8$  NS(5, -709.13261, 5.808;  $-\infty$ ) with  $\rho_\delta = 1.653$   
 NS(5, -0.334453434, 7/9; 0) with  $\rho_\delta = 9/7$   
 We could not find periodic points of period 3.  
 $q = 10$  NS(9, -0.17716996425, 9/11; -0.02106166)  
 NS(9, -0.620995042, 9/11; -0.5302298)  
 NS(9, -9110.600525, 9/11; -21236.960613) with  $\rho_\delta = 9/11$   
 NS(9, -29708498.34, 9/11; -153703407.949836) with  $\rho_\delta = 9/11$   
 We could not find stable periodic points of period 7.  
 $q = 12$  NS(7, -0.2792913671, 11/13; -0.1789) with  $\rho_\delta = 11/13$   
 NS(7, -23823237.5, 11/13; -1194427400) with  $\rho_\delta = 11/13$   
 NS(9, -31111814756509686, 11/13;  $-31793 \times 10^{10}$ ) with  $\rho_\delta = 11/13$   
 We could not find stable periodic points of period 5.

Thus it is very likely that these results support

**Conjecture H :** For  $q \neq 1, 2, 4$  there are geometrically convergent sequences

$$\lambda(m), \lambda(m+1), \lambda(m+2), \dots$$

with asymptotic ratio  $(q-1)/(q+1)$ , whose orbits have diameters such that  $\rho_\delta$  are equal to  $(q-1)/(q+1)$  or  $(q+1)/(q-1)$ .

### References

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Table 1  $\lambda(n)$  and  $\rho(n)$  for Newton iteration of  $x^q + (\lambda - 1)x - \lambda$

$q = 3$

$n$	$\lambda(n)$	$\delta(n)$	$\rho(n)$
2	0.36110 30805 2	$-3.289742 \times 10^{-2}$	
3	0.51065 04599 9	$-3.50883972 \times 10^{-3}$	
4	0.64180 42757 2	$-2.42792157 \times 10^{-3}$	0.8770051083129
5	0.74542 65254 3	$-1.40180586 \times 10^{-3}$	0.7900818526189
6	0.82270 63415 2	$-7.3382092 \times 10^{-4}$	0.7457840020485
7	0.87824 54652 5	$-3.6231445 \times 10^{-4}$	0.7186756716051
...			
33	0.99999 65839 25232 985	$-3.08945 \times 10^{-13}$	0.6666680287358
34	0.99999 77226 15581 208	$-1.37309 \times 10^{-13}$	0.6666675747132
35	0.99999 84817 43169 34703	$-6.102634 \times 10^{-14}$	0.66666727203200
36	0.99999 89878 28534 471	$-2.7123 \times 10^{-14}$	0.66666670702413
...			
3	0.26744 849	$-9.018 \times 10^{-5}$	
4	0.35689 6844	$-9.02511 \times 10^{-4}$	
5	0.39875 6368	$-1.31541 \times 10^{-4}$	0.467974223427297
6	0.41724 9849	$-4.4089 \times 10^{-5}$	0.441798645310098
7	0.42813 61148	$-1.56378 \times 10^{-5}$	0.588654229022648
...			
35	0.44598 71767 52265 33855	$-1.07137 \times 10^{-16}$	0.6666647553720277458
36	0.44598 72354 92575 02550 9	$-4.7616 \times 10^{-17}$	0.6666653892430334713
37	0.44598 72746 52731 45919 1	$-1.49643 \times 10^{-17}$	0.66666581504914995173
38	0.44598 73007 59480 18195 5	$-6.66506 \times 10^{-18}$	0.66666609892048386

$q = 5$

3	0.46346 9485	$-3.719969 \times 10^{-3}$	
4	0.53746 9858	$-4.112064 \times 10^{-3}$	
5	0.60508 4424	$-4.06828 \times 10^{-3}$	0.91370574578
6	0.66588 47865	$-3.7769341 \times 10^{-3}$	0.899219888507
7	0.71972 72770	$-3.3528685 \times 10^{-3}$	0.88556199808
8	0.76669 77085	$-2.8760334 \times 10^{-3}$	0.87236736383
...			

59	0.99999 68241 07403	$-1.471601 \times 10^{-9}$	0.8000009864352259
60	0.99999 74592 84819	$-1.0932 \times 10^{-10}$	0.8000007728241934
61	0.99999 79674 27149	$-8.12083 \times 10^{-10}$	0.8000006253370947
62	0.99999 83739 41267	$-6.03244 \times 10^{-10}$	0.8000004998599506

$q=7$

5	0.42312 5779	$-3.03754 \times 10^{-4}$	
6	0.43122 22136	$-1.801304 \times 10^{-4}$	
7	0.43623 54287	$-1.107137 \times 10^{-4}$	0.619187994
8	0.43948 89649	$-7.04182 \times 10^{-5}$	0.648991941
9	0.44169 22537	$-4.6579 \times 10^{-5}$	0.677198182
...			
52	0.44876 20004 6141	$-5.69932 \times 10^{-9}$	0.8856968078787
53	0.44876 27803 91479	$-4.743642 \times 10^{-9}$	0.8569930614147
54	0.44876 34488 02827	$-3.948136 \times 10^{-9}$	0.8570144712719
55	0.44876 40216 53287	$-3.285965 \times 10^{-9}$	0.8570328162651
...			
3	0.45820 702	$-2.34891 \times 10^{-3}$	
4	0.50885 7184	$-2.749562 \times 10^{-3}$	
5	0.55631 61443	$-2.9436375 \times 10^{-3}$	0.9369951951192
6	0.60077 45068	$-3.0076572 \times 10^{-3}$	0.93677489390765
7	0.64230 45144	$-2.976515 \times 10^{-3}$	0.934132641524977
...			
80	0.99999 35879 69	$-1.4176 \times 10^{-8}$	0.85714383927049130
81	0.99999 45039 69	$-1.1792 \times 10^{-8}$	0.85714491960852568
82	0.99999 52891 139	$-9.81 \times 10^{-9}$	0.85714339285489198
83	0.99999 59620 956	$-8.1606 \times 10^{-9}$	0.85714508733624454



Table 2  $\lambda (n)$  and  $\rho (n)$  for Halley iteration of  $x^q + (\lambda - 1)x - \lambda$

$n$	$\lambda (n)$	$\delta (n)$	$\rho (n)$
$q = 3$			
3	1.92678 79381 8	$-4.06252683 \times 10^{-3}$	
4	2.19168 17474 7	$-1.35657533 \times 10^{-3}$	
5	2.33984 12999 2	$-3.9016401 \times 10^{-4}$	0.5593167799848
6	2.41833 95056 9	$-1.0453124 \times 10^{-4}$	0.5298221037519
7	2.45876 36191 50	$-2.7048023 \times 10^{-5}$	0.5149686297091
...			
40	2.49999 99999 95151 29856 08105 72576	$-2.52910 \times 10^{-25}$	0.50000000001750327
41	2.49999 99999 97575 64928 71820 50648	$-6.322739 \times 10^{-26}$	0.500000001398812097
42	2.49999 99999 98787 82464 02015 82293 309	$-1.58069 \times 10^{-26}$	0.499999995806626788
43	2.49999 99999 99393 91232 01007 02742 41	$-3.95171 \times 10^{-27}$	0.500000002796092675
...			
5	1.58654 04026 18	$-8.035835 \times 10^{-4}$	
6	1.61576 23172 48	$-1.5485105 \times 10^{-5}$	
7	1.62913 60179 39	$-3.440801 \times 10^{-6}$	0.457659974041
8	1.63554 49315 05	$-8.11174 \times 10^{-7}$	0.4792176611454
...			
33	1.64177 42807 94355 59509 04108 20	$-4.5272 \times 10^{-22}$	0.4999999993562528
34	1.64177 42808 85856 66822 22906 22	$-1.13179 \times 10^{-22}$	0.4999999996781264
35	1.64177 42809 31607 20477 35046 346	$-2.82947 \times 10^{-23}$	0.4999999998390632
36	1.64177 42809 54482 47304 54301 686	$-7.0737 \times 10^{-24}$	0.4999999999195316
$q = 5$			
3	2.61686 822	$-2.36051 \times 10^{-2}$	
4	4.28192 897	$-2.26933 \times 10^{-2}$	
5	6.38844 1112	$-1.7253962 \times 10^{-2}$	1.265126296
6	8.61083 7053	$-1.100877 \times 10^{-2}$	1.05501216859
7	10.64030 59509	$-6.2167065 \times 10^{-3}$	0.91318961646
...			
47	16.64114 90921 32836 07734 477	$-5.417471 \times 10^{-17}$	0.6666666856857448
48	16.64114 93228 35276 96763 782	$-2.407764 \times 10^{-17}$	0.66666667934605204
49	16.64114 94766 36906 17794 329	$-1.070117 \times 10^{-17}$	0.66666667511959024
50	16.64114 95791 71326 51819 588	$-4.75607 \times 10^{-17}$	0.66666667230194903

51	16.64114 96475 27607 13023 788	$-2.11381 \times 10^{-17}$	0.66666667042352157
$q = 7$			
5	2.05839 513	$-6.9965 \times 10^{-4}$	
6	2.23228 5997	$-2.16289 \times 10^{-4}$	
7	2.34058 0864	$-8.0135 \times 10^{-5}$	0.6227749233
8	2.41109 6182	$-3.3293 \times 10^{-5}$	0.6511418311
9	2.45864 51793	$-1.50151 \times 10^{-5}$	0.67430734695
...			
47	2.57822 93579 37855 47716	$-1.71307 \times 10^{-15}$	0.74999 86614 578
48	2.57822 98253 70780 73004	$-9.6361 \times 10^{-16}$	0.74999 89960 937
49	2.57823 01759 45122 72559	$-5.4203 \times 10^{-16}$	0.74999 92459 1075
50	2.57823 04388 75681 25392 6	$-3.0489 \times 10^{-16}$	0.74999 94353 0284
51	2.57823 06360 73488 79308 54	$-1.715003 \times 10^{-16}$	0.74999 95764 7719
...			
3	5.85146 057	-0.174595	
4	47.29110 1	-1.3426689	
5	361.53335 62	-5.7101412	7.581302321
6	1696.707076	-14.127813	4.248867546
7	5435.72649 42	-24.3720025	2.800399201
8	13176.37436 53	-32.4094343	2.0702347341
9	25917.58168 89	-35.387315	1.6460130386
10	43469.59230 84	-33.175574	1.3775782917
11	64505.60010 947	-27.622842	1.1984956172
...			
82	220649.30705 45211 89604 96855 4	$-1.1493 \times 10^{-16}$	0.7500000005283192
83	220649.30714 89563 69876 62533 6	$-6.4648 \times 10^{-17}$	0.7500000003962393
84	220649.30721 97827 55108 43212 0	$-3.6364 \times 10^{-17}$	0.7500000002971797
85	220649.30727 29025 44048 07332 05	$-2.0455 \times 10^{-17}$	0.7500000002228839